

parallel-plate channel and a circular tube, respectively, plotted against the axial distance in the range $10^{-4} \leq Z \leq 10^{-1}$ for several different values of the dimensionless time τ . Starting from the inlet region, the local Nusselt numbers decrease continuously with both increasing time and axial location along the conduit until the conduction region is reached. In the conduction region, the Nusselt number remains invariant with the position but decreases with increasing time. Eventually, with increasing time, the local Nusselt numbers for both regions assume the well-known steady-state value.

REFERENCES

1. E. M. Sparrow and R. Siegel, Thermal entrance region of a circular tube under transient heating conditions, *Proceedings of the Third U.S. Congress for Applied Mechanics*, ASME, pp. 817-826 (1958).
2. R. Siegel and E. M. Sparrow, Transient heat transfer for laminar forced convection in the thermal entrance region of flat ducts, *J. Heat Transfer* **81**, 29-36 (1959).
3. R. Siegel, Heat transfer for laminar flow in ducts with arbitrary time variations in wall temperature, *J. Appl. Mech.* **27**, 241-249 (1960).
4. H. T. Lin and Y. P. Shih, Unsteady thermal entrance heat transfer of power-law fluids in pipes and plate slits, *Int. J. Heat Mass Transfer* **24**, 1531-1539 (1981).
5. S. C. Chen, N. K. Anand and D. R. Tree, Analysis of transient laminar convective heat transfer inside a circular duct, *J. Heat Transfer* **105**, 922-924 (1983).
6. R. M. Cotta and M. N. Özişik, Transient forced convection in laminar channel flow with timewise variations of wall temperature, ASME Paper No. 85-WA/HT-72 (1985).
7. M. N. Özişik and R. L. Murray, On the solution of linear diffusion problems with variable boundary condition parameters, *J. Heat Transfer* **96**, 48-51 (1974).
8. M. D. Mikhailov and M. N. Özişik, *Unified Analysis and Solutions of Heat and Mass Diffusion*. Wiley, New York (1984).
9. M. D. Mikhailov, M. N. Özişik and N. L. Vulchanov, Transient heat diffusion in one-dimensional composite media and automatic solution of the eigenvalue problem, *Int. J. Heat Mass Transfer* **26**, 1131-1141 (1983).
10. M. D. Mikhailov and N. L. Vulchanov, A computational procedure for Sturm-Liouville problems, *J. Comp. Phys.* **50**, 323-336 (1983).

Natural convection heat transfer in enclosures with an off-center partition

TATSUO NISHIMURA, MITSUHIRO SHIRAISHI and YUJI KAWAMURA

Department of Chemical Engineering, Hiroshima University, Higashi-Hiroshima 724, Japan

(Received 1 December 1986 and in final form 19 January 1987)

1. INTRODUCTION

NATURAL convection through externally heated and cooled enclosures is of interest in solar collector applications, in the estimation of heat losses from double-pane windows, and in the calculation of heat losses from rooms. Numerous experimental and numerical computational studies have been reported explaining the heat transfer mechanism and presenting correlations for heat transfer rates for such systems. Excellent reviews [1, 2] are available and there is no need to repeat them here.

The problem of primary interest in the literature [1, 2] is that of an enclosure with no partitions. However, in practical cases, a vertical partition is inserted into the enclosure to reduce heat losses by natural convection and thermal radiation. Reported studies of natural convection in a partitioned enclosure are limited. Duxbury [3] reported experiments with air-filled enclosures containing a central partition as shown in Fig. 1. Nakamura *et al.* [4] performed computational and experimental studies including the effect of thermal radiation for the same configuration as that of Duxbury. The present authors [5] proposed a boundary layer solution for this system and confirmed its validity by experiments. Tong and Gerner [6] reported the effect of partition position on the heat transfer rate by numerical computation and concluded that a central partition corresponding to $W'/W = 0.5$ produces the greatest reduction in heat transfer.

This study is an extension of the previous study [5]. We examine the limitations of the boundary layer approximation for various positions of the partition. We show that even if the partition deviates from the center of the enclosure, the heat transfer rate is identical with that for the partition in the central position. This does not appear to have been studied previously.

2. LIMITATIONS OF THE BOUNDARY LAYER APPROXIMATION

We [5] previously indicated that a thermal boundary layer with a constant thickness is developed along the partition at high Rayleigh numbers for the enclosure with a central

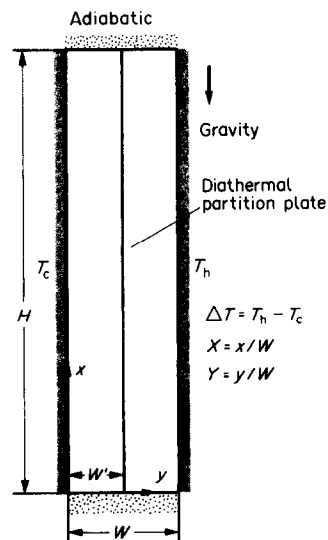


FIG. 1. Schematic diagram of an enclosure divided by a vertical partition.

NOMENCLATURE

g	gravitational acceleration [m s ⁻²]	W'	distance between the partition and the cold wall [m]
H	height of the enclosure [m]	W_{\min}	minimum width of cell satisfying the boundary layer approximation [m].
h	heat transfer coefficient, $q/(T_h - T_c)$ [W m ⁻² K ⁻¹]		
L	depth of the enclosure [m]		
Nu	average Nusselt number, $h/W\lambda$ [—]		
Pr	Prandtl number [—]		
q	heat flux per unit area [W m ⁻²]		
Ra	Rayleigh number, $g\beta(T_h - T_c)W^3/(\alpha\nu)$ [—]		
T	temperature [K]		
T_c	temperature at the cold wall [K]		
T_h	temperature at the hot wall [K]		
W	width of the enclosure [m]		

Greek symbols

α	thermal diffusivity [m ² s ⁻¹]
β	volumetric expansion coefficient [K ⁻¹]
δ	thermal boundary layer thickness defined in Fig. 2 [m]
λ	thermal conductivity [W m ⁻¹ K ⁻¹]
ν	kinematic viscosity [m ² s ⁻¹].

partition and presented the boundary layer solution predicting the heat transfer rate through the enclosure.

The temperature distribution proposed by the boundary layer model is shown in Fig. 2. The thickness of the boundary layer is defined as the distance within which the temperature near the partition reaches the core temperature as indicated in Fig. 2. The boundary layer thickness is given by

$$\delta/W = 2.64Ra^{-1/4}(H/W)^{1/4}. \tag{1}$$

In this study it is shown that the same relation is applied for certain values of the boundary layer thickness regardless of the position of the partition. That is, it may be considered from the boundary layer model that the heat transfer rate is independent of the position of the partition if the boundary layer thickness is less than the half-width of each cell constructed by the partition, i.e. $\delta < W'/2$ and $\delta < (W - W')/2$. Thus the minimum width of a cell satisfying the boundary

layer approximation is derived by using equation (1)

$$W_{\min}/W = 5.28Ra^{-1/4}(H/W)^{1/4}. \tag{2}$$

The validity of equation (2) is confirmed by experimental measurement and numerical computation as subsequently described.

3. EXPERIMENT

Since the experimental equipment and procedure have already been described in detail [7], they will be reviewed here only briefly. Two kinds of enclosure were used. The height and the length were fixed ($H = 300$ mm and $L = 200$ mm), and the width was variable ($W = 30$ and 75 mm). The vertical partition was made of aluminum foil, $15 \mu\text{m}$ in thickness, and its position was changeable ($W'/W = 0.166$ – 0.75). The working fluid was water ($Pr = 6$). The experiments were carried out in the range $10^6 < Ra < 10^8$.

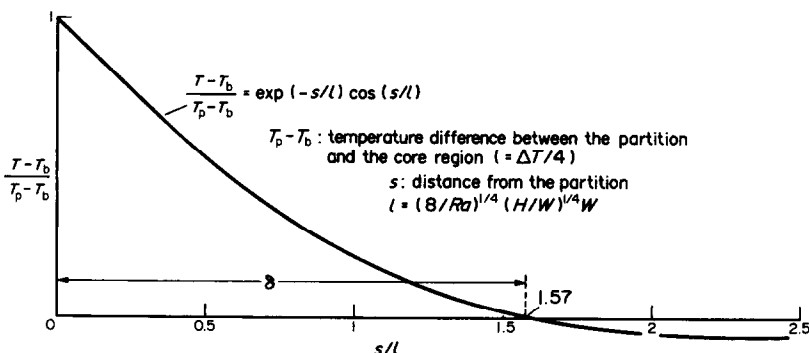


FIG. 2. Temperature profile near the partition proposed by the boundary layer model.

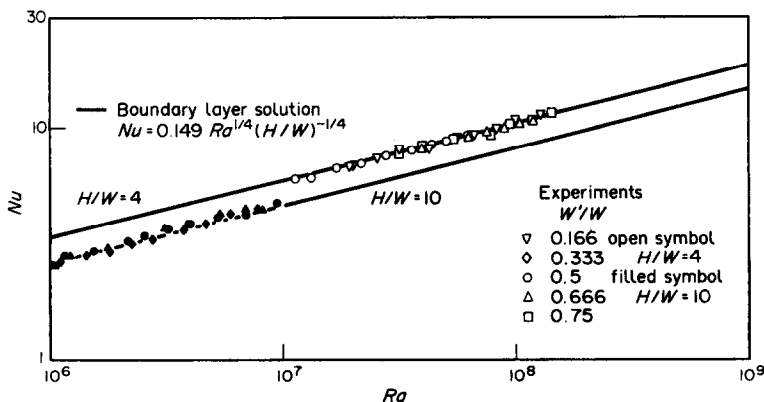


FIG. 3. Experimental Nusselt number vs Rayleigh number for various positions of the partition.

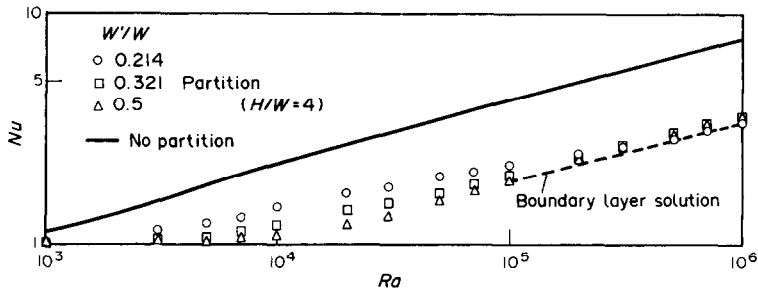


FIG. 4. Experimental Nusselt number vs Rayleigh number for various positions of the partition.

The results of the experiments are shown in Fig. 3 in which the Nusselt number is expressed as a function of the Rayleigh number. The Nusselt numbers for a given Rayleigh number are independent of various values of W'/W for the values of $H/W = 4$ and 10. Also, all of experimental data agree well with the boundary layer solution presented previously for the case of a central partition. Under these experimental conditions, the width of each cell constructed by the partition is larger than the minimum width calculated using equation (2), and thus the premise described in Section 2 is valid.

4. NUMERICAL COMPUTATION

A theoretical investigation was carried out to confirm the validity of equation (2) using a two-dimensional Galerkin finite element method. The solution technique is the same as that previously used [8]. The equations utilized were the Navier-Stokes and the energy transport equations. The computation was performed in the range $10^3 < Ra < 10^6$ for $Pr = 6$ and $H/W = 4$. Solutions were obtained for different positions of the partition from the center of the enclosure to the cold wall ($W'/W < 0.5$). Because the experimental Nusselt numbers are essentially identical whether the position of the partition is located at the cold or hot wall side as shown in Fig. 3.

Figure 4 shows the calculated results of the Nusselt number as a function of the Rayleigh number. In this case, Nusselt numbers for various values of W'/W deviate from the boundary layer solution, particularly for Rayleigh numbers less than 10^5 . Equation (2) predicts that the boundary layer approximation is satisfied for W'/W larger than 0.418 at $Ra = 10^5$. This agrees with the results of the numerical computation at $Ra = 10^5$.

Thus the validity of equation (2) is considered to be confirmed both by experimental measurement and numerical computation as indicated above.

Acknowledgements—The authors acknowledge with thanks the assistance of Mr Masayuki Yamamoto with the experi-

ments. This work was supported in part by a Grant-in-Aid for Scientific Research (No. 61750891) from the Ministry of Education, Science and Culture of Japan. We are also very grateful for the revision in the English text by Professor W. Hayduk of the University of Ottawa.

REFERENCES

1. I. Catton, Natural convection in enclosures, *Proc. 6th Int. Heat Transfer Conference*, Vol. 6, pp. 13–43 (1978).
2. S. W. Churchill, Free convection in layers and enclosures. In *Heat Exchanger Design Handbook*, Chap. 2.5.8. Hemisphere, Washington, DC (1983).
3. D. Duxbury, An interferometric study of natural convection in enclosed plane air layers with complete and partial central vertical divisions, Ph.D. Thesis, University of Salford (1979).
4. H. Nakamura, Y. Asako and T. Hirata, Natural convection and thermal radiation in enclosures with partition plate, *Trans. JSME Ser. B*, **50**, 2647–2654 (1984).
5. T. Nishimura, M. Shiraishi and Y. Kawamura, Analysis of natural convection heat transfer in enclosures divided by a vertical partition plate, *Proc. Int. Symp. Heat Transfer*, Beijing, Paper No. 85-ISHT-1-6 (1985).
6. T. W. Tong and F. M. Gerner, Natural convection in partitioned air-filled rectangular enclosures, *Int. Commun. Heat Mass Transfer* **13**, 99–108 (1986).
7. T. Nishimura, T. Takumi, Y. Kawamura and H. Ozoe, Experiments of natural convection heat transfer in rectangular enclosure partially filled with particles, *Kagaku Kogaku Ronbunshu* **11**, 405–411 (1985); *Heat Transfer—Jap. Res.* **15**, 62–76 (1986).
8. T. Nishimura, T. Takumi, M. Shiraishi, Y. Kawamura and H. Ozoe, Numerical analysis of natural convection in a rectangular enclosure horizontally divided into fluid and porous regions, *Int. J. Heat Mass Transfer* **29**, 889–898 (1986).

A generalized correlation for thermal design data of heat-pipe heat exchangers

SHOU-SHING HSIEH

Department of Mechanical Engineering, National Sun Yat-Sen University,
Kaohsiung, Taiwan 80424, Republic of China

(Received 16 May 1986 and in final form 30 January 1987)

INTRODUCTION

HEAT EXCHANGERS made of heat pipes have attracted much attention in the application of economic devices for the

recovery of waste heat energy [1–3]. Although the characteristics of thermal performance of a single heat pipe have been extensively studied and clearly understood during the past 20 years, the study of the overall performance of heat